

Spectral Zeta Functions and Scalar Curvature for Noncommutative Tori

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Sharif University, Tehran, Dec. 2012

- ▶ A. Connes and P. Tretkoff, *The Gauss-Bonnet Theorem for the noncommutative two torus*, Sept. 2009.
- ▶ A. Connes and H. Moscovici, *Modular curvature for noncommutative two-tori*, Oct. 2011.
- ▶ F. Fathizadeh and M. Khalkhali, *The Gauss-Bonnet Theorem for noncommutative two tori with a general conformal structure*, May 2010.
- ▶ F. Fathizadeh and M. Khalkhali, *Scalar Curvature for the Noncommutative Two Torus*, Oct. 2011.

Laplace spectrum; commutative background

- ▶ (M, g) = closed Riemannian manifold. [Laplacian on forms](#)

$$\Delta = (d + d^*)^2 : \Omega^p(M) \rightarrow \Omega^p(M),$$

has pure point spectrum:

$$0 \leq \lambda_1 \leq \lambda_2 \leq \dots \rightarrow \infty$$

- ▶ Fact: [Dimension](#), [volume](#), [total scalar curvature](#), [Betti numbers](#), and hence the [Euler characteristic](#) of M are fully determined by the spectrum of Δ (on all p -forms).

Method of proof: bring in the heat kernel

- ▶ Heat equation for functions: $\partial_t + \Delta = 0$
- ▶ $k(t, x, y)$ = kernel of $e^{-t\Delta}$. Asymptotic expansion near $t = 0$:

$$k(t, x, x) \sim \frac{1}{(4\pi t)^{m/2}} (a_0(x, \Delta) + a_1(x, \Delta)t + a_2(x, \Delta)t^2 + \dots)$$

- ▶ $a_i(x, \Delta)$, Seeley-De Witt-Gilkey coefficients.

- Theorem: $a_i(x, \Delta)$ are universal polynomials in curvature tensor R and its covariant derivatives:

$$a_0(x, \Delta) = 1$$

$$a_1(x, \Delta) = \frac{1}{6}S(x) \quad \text{scalar curvature}$$

$$a_2(x, \Delta) = \frac{1}{360}(2|R(x)|^2 - 2|\text{Ric}(x)|^2 + 5|S(x)|^2)$$

$$a_3(x, \Delta) = \dots\dots\dots$$

Heat trace asymptotics

Compute $\text{Trace}(e^{-t\Delta})$ in two ways:

Spectral Sum = Geometric Sum.

$$\sum e^{-t\lambda_i} = \int_M k(t, x, x) d\text{vol}_x \sim (4\pi t)^{-\frac{m}{2}} \sum_{j=0}^{\infty} a_j t^j \quad (t \rightarrow 0).$$

Hence

$$a_j = \int_M a_j(x, \Delta) d\text{vol}_x,$$

are manifestly spectral invariants:

$$a_0 = \int_M d\text{vol}_x = \text{Vol}(M), \quad \implies \text{Weyl's law}$$

$$a_1 = \frac{1}{6} \int_M S(x) d\text{vol}_x, \quad \text{total scalar curvature}$$

Tauberian theory and $a_0 = 1$, implies [Weyl's law](#):

$$N(\lambda) \sim \frac{\text{Vol}(M)}{(4\pi)^{m/2} \Gamma(1 + m/2)} \lambda^{m/2} \quad \lambda \rightarrow \infty,$$

where

$$N(\lambda) = \#\{\lambda_i \leq \lambda\}$$

is the eigenvalue counting function.

Simplest example: flat tori

- ▶ $\Gamma \subset \mathbb{R}^m$ a cocompact lattice; $M = \mathbb{R}^m/\Gamma$

$$\text{spec}(\Delta) = \{4\pi^2 \|\gamma^*\|^2; \gamma^* \in \Gamma^*\}$$

- ▶ Then:

$$K(t, x, y) = \frac{1}{(4\pi t)^{m/2}} \sum_{\gamma \in \Gamma} e^{-4\pi^2 \|x-y+\gamma\|^2/4t}$$

- ▶ Poisson summation formula \implies

$$\sum_{\gamma^* \in \Gamma^*} e^{-4\pi^2 \|\gamma^*\|^2 t} = \frac{\text{Vol}(M)}{(4\pi t)^{m/2}} \sum_{\gamma \in \Gamma} e^{-4\pi^2 \|\gamma\|^2/4t}$$

- ▶ And from this we obtain the asymptotic expansion of the heat trace near $t = 0$

$$\text{Tr} e^{-t\Delta} \sim \frac{\text{Vol}(M)}{(4\pi t)^{m/2}} \quad (t \rightarrow 0)$$

Application 1: heat equation proof of the Atiyah-Singer index theorem

- ▶ Dirac operator

$$D : C^\infty(S_+) \rightarrow C^\infty(S_-)$$

McKean-Singer formula:

$$\text{Index}(D) = \text{Tr}(e^{-tD^*D}) - \text{Tr}(e^{-tDD^*}), \quad \forall t > 0$$

Heat trace asymptotics \implies

$$\text{Index}(D) = \int_M a_n(x) dx,$$

where $a_n(x) = a_n^+(x) - a_n^-(x)$, $m = 2n$, can be explicitly computed and recovers the A-S integrand (The simplest proof is due to Getzler).

Application 2: meromorphic extension of spectral zeta functions

$$\zeta_{\Delta}(s) := \sum_{\lambda_j \neq 0} \lambda_j^{-s}, \quad \operatorname{Re}(s) > \frac{m}{2}$$

Mellin transform + asymptotic expansion:

$$\lambda^{-s} = \frac{1}{\Gamma(s)} \int_0^{\infty} e^{-t} t^{s-1} dt \quad \operatorname{Re}(s) > 0$$

$$\begin{aligned} \zeta_{\Delta}(s) &= \frac{1}{\Gamma(s)} \int_0^{\infty} (\operatorname{Trace}(e^{-t\Delta}) - \dim \operatorname{Ker} \Delta) t^{s-1} dt \\ &= \frac{1}{\Gamma(s)} \left\{ \int_0^c \dots + \int_c^{\infty} \dots \right\} \end{aligned}$$

The second term defines an entire function, while the first term has a meromorphic extension to \mathbb{C} with **simple poles** within the set

$$\frac{m}{2} - j, \quad j = 0, 1, \dots$$

Also: 0 is always a regular point.

Simplest example: For $M = S^1$ with round metric, we have

$$\zeta_{\Delta}(s) = 2\zeta(2s) \quad \text{Riemann zeta function}$$

Scalar curvature

The spectral invariants a_j in the heat asymptotic expansion

$$\text{Trace}(e^{-t\Delta}) \sim (4\pi t)^{-\frac{m}{2}} \sum_{j=0}^{\infty} a_j t^j \quad (t \rightarrow 0)$$

are related to residues of spectral zeta function by

$$\text{Res}_{s=\alpha} \zeta_{\Delta}(s) = (4\pi)^{-\frac{m}{2}} \frac{a_{\frac{m}{2}-\alpha}}{\Gamma(\alpha)}, \quad \alpha = \frac{m}{2} - j > 0$$

Focusing on subleading pole $s = \frac{m}{2} - 1$ and using $a_1 = \frac{1}{6} \int_M S(x) d\text{vol}_x$, we obtain a formula for scalar curvature density as follows:

Let $\zeta_f(s) := \text{Tr}(f\Delta^{-s})$, $f \in C^\infty(M)$.

$$\text{Res } \zeta_f(s)|_{s=\frac{m}{2}-1} = \frac{(4\pi)^{-m/2}}{\Gamma(m/2-1)} \int_M fS(x) d\text{vol}_x, \quad m \geq 3$$

$$\zeta_f(s)|_{s=0} = \frac{1}{4\pi} \int_M fS(x) d\text{vol}_x - \text{Tr}(fP) \quad m = 2$$

$\log \det(\Delta) = -\zeta'(0)$, Ray-Singer regularized determinant

Noncommutative Geometry: Spectral Triples $(\mathcal{A}, \mathcal{H}, D)$

- ▶ \mathcal{A} = involutive unital algebra, \mathcal{H} = Hilbert space,

$$\pi : \mathcal{A} \rightarrow \mathcal{L}(\mathcal{H}), \quad D : \mathcal{H} \rightarrow \mathcal{H}$$

D has compact resolvent and all commutators $[D, \pi(a)]$ are bounded.

- ▶ An asymptotic expansion holds

$$\text{Trace}(e^{-tD^2}) \sim \sum a_\alpha t^\alpha \quad (t \rightarrow 0)$$

- ▶ Let $\Delta = D^2$. Spectral zeta function

$$\zeta_D(s) = \text{Tr}(|D|^{-s}) = \text{Tr}(\Delta^{-s/2}), \quad \text{Re}(s) \gg 0.$$

Curved noncommutative tori A_θ

$A_\theta = C(\mathbb{T}_\theta^2) =$ universal C^* -algebra generated by unitaries U and V

$$VU = e^{2\pi i\theta} UV.$$

$$A_\theta^\infty = C^\infty(\mathbb{T}_\theta^2) = \left\{ \sum_{m,n \in \mathbb{Z}} a_{m,n} U^m V^n : (a_{m,n}) \text{ Schwartz class} \right\}.$$

► Differential operators on A_θ

$$\delta_1, \delta_2 : A_\theta^\infty \rightarrow A_\theta^\infty,$$

Infinitesimal generators of the action

$$\alpha_s(U^m V^n) = e^{is \cdot (m,n)} U^m V^n \quad s \in \mathbb{R}^2.$$

Analogues of $\frac{1}{i} \frac{\partial}{\partial x}$, $\frac{1}{i} \frac{\partial}{\partial y}$ on 2-torus.

► Canonical trace $\mathfrak{t} : A_\theta \rightarrow \mathbb{C}$ on smooth elements:

$$\mathfrak{t}\left(\sum_{m,n \in \mathbb{Z}} a_{m,n} U^m V^n\right) = a_{0,0}.$$

Complex structures on A_θ

► Let $\mathcal{H}_0 = L^2(A_\theta)$ = GNS completion of A_θ w.r.t. \mathfrak{t} .

► Fix $\tau = \tau_1 + i\tau_2$, $\tau_2 = \Im(\tau) > 0$, and define

$$\partial := \delta_1 + \tau\delta_2, \quad \partial^* := \delta_1 + \bar{\tau}\delta_2.$$

► Hilbert space of $(1, 0)$ -forms:

$\mathcal{H}^{(1,0)}$:= completion of finite sums $\sum a\partial b$, $a, b \in A_\theta^\infty$, w.r.t.

$$\langle a\partial b, a'\partial b' \rangle := \mathfrak{t}((a'\partial b')^* a\partial b).$$

► Flat Dolbeault Laplacian: $\partial^*\partial = \delta_1^2 + 2\tau_1\delta_1\delta_2 + |\tau|^2\delta_2^2$.

Conformal perturbation of the metric

- ▶ Fix $h = h^* \in A_\theta^\infty$. Replace the volume form \mathfrak{t} by $\varphi : A_\theta \rightarrow \mathbb{C}$,

$$\varphi(a) := \mathfrak{t}(ae^{-h}), \quad a \in A_\theta.$$

- ▶ It is a twisted trace (in fact a KMS state)

$$\varphi(ab) = \varphi(b\Delta(a)), \quad \forall a, b \in A_\theta.$$

where

$$\Delta(x) = e^{-h}xe^h,$$

is the modular automorphism of a von Neumann factor-has no commutative counterpart.

- ▶ Warning: \triangle and Δ are very different operators!

Connes-Tretkoff spectral triple

- ▶ Hilbert space $\mathcal{H}_\varphi :=$ GNS completion of A_θ w.r.t. $\langle \cdot, \cdot \rangle_\varphi$,

$$\langle a, b \rangle_\varphi := \varphi(b^* a), \quad a, b \in A_\theta$$

- ▶ View $\partial_\varphi = \partial = \delta_1 + \tau\delta_2 : \mathcal{H}_\varphi \rightarrow \mathcal{H}^{(1,0)}$. and let

$$\partial_\varphi^* : \mathcal{H}^{(1,0)} \rightarrow \mathcal{H}_\varphi$$

$$\mathcal{H} = \mathcal{H}_\varphi \oplus \mathcal{H}^{(1,0)},$$

$$D = \begin{pmatrix} 0 & \partial_\varphi^* \\ \partial_\varphi & 0 \end{pmatrix} : \mathcal{H} \rightarrow \mathcal{H}.$$

Full perturbed Laplacian:

$$\Delta := D^2 = \begin{pmatrix} \partial_\varphi^* \partial_\varphi & 0 \\ 0 & \partial_\varphi \partial_\varphi^* \end{pmatrix} : \mathcal{H} \rightarrow \mathcal{H}.$$

Lemma: $\partial_\varphi^* \partial_\varphi : \mathcal{H}_\varphi \rightarrow \mathcal{H}_\varphi$, and $\partial_\varphi \partial_\varphi^* : \mathcal{H}^{(1,0)} \rightarrow \mathcal{H}^{(1,0)}$ are anti-unitarily equivalent to

$$\begin{aligned} k \partial^* \partial k &: \mathcal{H}_0 \rightarrow \mathcal{H}_0, \\ \partial^* k^2 \partial &: \mathcal{H}^{(1,0)} \rightarrow \mathcal{H}^{(1,0)}, \end{aligned}$$

where $k = e^{h/2}$.

Scalar curvature for A_θ

- ▶ The scalar curvature of the curved nc torus $(\mathbb{T}_\theta^2, \tau, k)$ is the unique element $R \in A_\theta^\infty$ satisfying

$$\text{Trace}(a\Delta^{-s})|_{s=0} + \text{Trace}(aP) = \text{t}(aR), \quad \forall a \in A_\theta^\infty,$$

where P is the projection onto the kernel of Δ .

- ▶ In practice this is done by finding an asymptotic expansion for the kernel of the operator $e^{-t\Delta}$, using Connes' [pseudodifferential calculus](#) for nc tori. A good pseudo diff calculus for general nc spaces is still illusive.

Connes' pseudodifferential calculus

- ▶ Symbols: smooth maps $\rho : \mathbb{R}^2 \rightarrow A_\theta^\infty$. Ψ DO's: $P_\rho : A_\theta^\infty \rightarrow A_\theta^\infty$,

$$P_\rho(a) = (2\pi)^{-2} \int_{\mathbb{R}^2} \int_{\mathbb{R}^2} e^{-is \cdot \xi} \rho(\xi) \alpha_s(a) ds d\xi.$$

Even for polynomial symbols these integrals are badly divergent; make sense as Fourier integral operators.

- ▶ For example:

$$\rho(\xi_1, \xi_2) = \sum a_{ij} \xi_1^i \xi_2^j, \quad a_{ij} \in A_\theta^\infty \quad \Rightarrow \quad P_\rho = \sum a_{ij} \delta_1^i \delta_2^j.$$

- ▶ Multiplication of symbol.

$$\sigma(PQ) \sim \sum_{\ell_1, \ell_2 \geq 0} \frac{1}{\ell_1! \ell_2!} \partial_1^{\ell_1} \partial_2^{\ell_2} (\rho(\xi)) \delta_1^{\ell_1} \delta_2^{\ell_2} (\rho'(\xi)).$$

Local expression for the scalar curvature

- ▶ Cauchy integral formula:

$$e^{-t\Delta} = \frac{1}{2\pi i} \int_C e^{-t\lambda} (\Delta - \lambda)^{-1} d\lambda.$$

- ▶ $B_\lambda \approx (\Delta - \lambda)^{-1}$:

$$\sigma(B_\lambda) \sim b_0(\xi, \lambda) + b_1(\xi, \lambda) + b_2(\xi, \lambda) + \dots,$$

each $b_j(\xi, \lambda)$ is a symbol of order $-2 - j$, and

$$\sigma(B_\lambda(\Delta - \lambda)) \sim 1.$$

(Note: λ is considered of order 2.)

Proposition: The scalar curvature of the spectral triple attached to (A_θ, τ, k) is equal to

$$\frac{1}{2\pi i} \int_{\mathbb{R}^2} \int_C e^{-\lambda} b_2(\xi, \lambda) d\lambda d\xi,$$

where b_2 is defined as above.

The computations for $k\partial^*\partial k$

- ▶ The symbol of $k\partial^*\partial k$ is equal to

$$a_2(\xi) + a_1(\xi) + a_0(\xi)$$

where

$$a_2(\xi) = \xi_1^2 k^2 + |\tau|^2 \xi_2^2 k^2 + 2\tau_1 \xi_1 \xi_2 k^2,$$

$$a_1(\xi) = 2\xi_1 k \delta_1(k) + 2|\tau|^2 \xi_2 k \delta_2(k) + 2\tau_1 \xi_1 k \delta_2(k) + 2\tau_1 \xi_2 k \delta_1(k),$$

$$a_0(\xi) = k \delta_1^2(k) + |\tau|^2 k \delta_2^2(k) + 2\tau_1 k \delta_1 \delta_2(k).$$

- ▶ The equation

$$(b_0 + b_1 + b_2 + \dots)((a_2 + 1) + a_1 + a_0) \sim 1,$$

has a solution with each b_j a symbol of order $-2 - j$.

$$b_0 = (a_2 + 1)^{-1} = (\xi_1^2 k^2 + |\tau|^2 \xi_2^2 k^2 + 2\tau_1 \xi_1 \xi_2 k^2 + 1)^{-1},$$

$$b_1 = -(b_0 a_1 b_0 + \partial_1(b_0) \delta_1(a_2) b_0 + \partial_2(b_0) \delta_2(a_2) b_0),$$

$$b_2 = -(b_0 a_0 b_0 + b_1 a_1 b_0 + \partial_1(b_0) \delta_1(a_1) b_0 + \partial_2(b_0) \delta_2(a_1) b_0 + \\ \partial_1(b_1) \delta_1(a_2) b_0 + \partial_2(b_1) \delta_2(a_2) b_0 + (1/2) \partial_{11}(b_0) \delta_1^2(a_2) b_0 + \\ (1/2) \partial_{22}(b_0) \delta_2^2(a_2) b_0 + \partial_{12}(b_0) \delta_{12}(a_2) b_0)$$

$$= 5\xi_1^2 b_0^2 k^3 \delta_1^2(k) b_0 + 2\xi_1^2 b_0 k \delta_1(k) b_0 \delta_1(k) b_0 k \\ + \text{about 800 terms.}$$

To perform the \mathbb{R}^2 integration and simplify, need a rearrangement lemma (Connes)

The computation of $\int_0^\infty \bullet r dr$ of these terms is achieved by:

For all $\rho_j \in A_\theta^\infty$ and $m_j > 0$ one has

$$\begin{aligned} & \int_0^\infty (k^2 u + 1)^{-m_0} \rho_1(k^2 u + 1)^{-m_1} \cdots \rho_\ell(k^2 u + 1)^{-m_\ell} u^{\sum m_j - 2} du \\ &= k^{-2(\sum m_j - 1)} F_{m_0, m_1, \dots, m_\ell}(\Delta_{(1)}, \Delta_{(2)}, \dots, \Delta_{(\ell)})(\rho_1 \rho_2 \cdots \rho_\ell), \end{aligned}$$

where

$$F_{m_0, m_1, \dots, m_\ell}(u_1, u_2, \dots, u_\ell) = \int_0^\infty (u + 1)^{-m} \prod_1^\ell (u \prod_1^j u_h + 1)^{-m_j} u^{\sum m_j - 2}$$

and $\Delta_{(i)}$ signifies that Δ acts on the i -th factor.

Final formula for the scalar curvature (Connes-Moscovici, Fathizadeh-K, Oct. 2011)

Theorem: The scalar curvature of (A_θ, τ, k) , up to an overall factor of $\frac{-\pi}{\tau_2}$, is equal to

$$\begin{aligned} & R_1(\log \Delta)(\Delta_0(\log k)) + \\ & R_2(\log \Delta_{(1)}, \log \Delta_{(2)}) \left(\delta_1(\log k)^2 + |\tau|^2 \delta_2(\log k)^2 + \tau_1 \{ \delta_1(\log k), \delta_2(\log k) \} \right) + \\ & iW(\log \Delta_{(1)}, \log \Delta_{(2)}) \left(\tau_2 [\delta_1(\log k), \delta_2(\log k)] \right) \end{aligned}$$

where

$$R_1(x) = -\frac{\frac{1}{2} - \frac{\sinh(x/2)}{x}}{\sinh^2(x/4)},$$

$$R_2(s, t) = (1 + \cosh((s + t)/2)) \times \frac{-t(s + t) \cosh s + s(s + t) \cosh t - (s - t)(s + t + \sinh s + \sinh t - \sinh(s + t))}{st(s + t) \sinh(s/2) \sinh(t/2) \sinh^2((s + t)/2)},$$

$$W(s, t) = -\frac{(-s - t + t \cosh s + s \cosh t + \sinh s + \sinh t - \sinh(s + t))}{st \sinh(s/2) \sinh(t/2) \sinh((s + t)/2)}.$$

The limiting case

In the commutative case, the above modular curvature reduces to a constant multiple of the [formula of Gauss](#):

$$\frac{1}{\tau_2} \delta_1^2(\log k) + \frac{|\tau|^2}{\tau_2} \delta_2^2(\log k) + 2 \frac{\tau_1}{\tau_2} \delta_1 \delta_2(\log k).$$

First application: Ray-Singer determinant and conformal anomaly (Connes-Moscovici)

Recall: $\log \text{Det}'(\Delta) = -\zeta'_{\Delta}(0)$, where Δ is the perturbed Laplacian on \mathbb{T}_{θ}^2 . One has the following *conformal variation formula*. Let $\nabla_i = \log \Delta$ which acts on the i -th factor of products.

Lemma

The log-determinant of the perturbed Laplacian Δ on \mathbb{T}_{θ}^2 is given by

$$\log \text{Det}'(\Delta) = \log \text{Det}' \Delta_0 + \log \varphi(1) - \frac{\pi}{12\tau_2} \varphi_0(h\Delta_0 h) - \frac{\pi}{4\tau_2} \varphi_0(K_2(\nabla_1)(\square_{\mathbb{R}}(h))),$$

Analogue of Osgood-Phillips-Sarnak functional on the space of selfadjoint elements of A_θ^∞ :

$$F(h) := -\log \text{Det}'(\Delta) + \log \varphi(1) = -\log \text{Det}'\left(e^{\frac{h}{2}} \Delta_0 e^{\frac{h}{2}}\right) + \log \varphi_0(e^{-h}).$$

Since $\Delta_{h+c} = e^c \Delta_h$ for any $c \in \mathbb{R}$, one has

$$\zeta_{\Delta_{h+c}}(z) = e^{-cz} \zeta_{\Delta_h}(z).$$

Therefore

$$\begin{aligned} F(h+c) &= \zeta'_{\Delta_{h+c}}(0) + \log \varphi_0(e^{-h-c}) \\ &= -c \zeta_{\Delta_h}(0) + \zeta'_{\Delta_h}(0) + \log \varphi_0(e^{-c}) + \log \varphi_0(e^{-h}) \\ &= F(h). \end{aligned}$$

Theorem

The functional $F(h)$ has the expression

$$F(h) = -(2 \log 2\pi + \log(|\eta(\tau)|^4)) + \frac{\pi}{4\tau_2} \varphi_0 \left(\left(\kappa_2 - \frac{1}{3} \right) (\nabla_1)(\square_{\mathfrak{R}}(h)) \right).$$

One has $F(h) \geq F(0)$ for all h and equality holds if and only if h is a scalar.

Second application: the Gauss-Bonnet theorem for A_θ

- ▶ Heat trace asymptotic expansion relates geometry to topology, thanks to MacKean-Singer formula:

$$\sum_{p=0}^m (-1)^p \text{Tr}(e^{-t\Delta_p}) = \chi(M) \quad \forall t > 0$$

- ▶ This gives the spectral formulation of the Gauss-Bonnet theorem:

$$\zeta(0) + 1 = \frac{1}{12\pi} \int_{\Sigma} R \, dv = \frac{1}{6} \chi(\Sigma)$$

Theorem (Connes-Tretkoff; Fathizadeh-K.): Let $\theta \in \mathbb{R}$, $\tau \in \mathbb{C} \setminus \mathbb{R}$, $k \in A_\theta^\infty$ be a positive invertible element. Then

$$\text{Trace}(\Delta^{-s})|_{s=0} + 2 = \text{t}(R) = 0,$$

where Δ is the Laplacian and R is the scalar curvature of the spectral triple attached to (A_θ, τ, k) .

The geometry in noncommutative geometry

- ▶ Geometry starts with **metric** and **curvature**. While there are a good number of 'soft' topological tools in NCG, like cyclic cohomology, K and KK-theory, and index theory, a truly noncommutative theory of curvature is still illusive. The situation is better with **scalar curvature**, but computations are quite tough at the moment.
- ▶ Metric aspects of NCG are informed by **Spectral Geometry**. Spectral invariants are the only means by which we can formulate metric ideas of NCG.