

Spectral Zeta Functions, Scalar Curvature, and Einstein-Hilbert action for Noncommutative Four Tori

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Related Works on \mathbb{T}_θ^2

A. Connes, P. Tretkoff, *The Gauss-Bonnet Theorem for the noncommutative two torus* (2009).

F. Fathizadeh, M. Khalkhali, *The Gauss-Bonnet Theorem for Noncommutative Two Tori With a General Conformal Structure* (2010), Journal of Noncommutative Geometry.

F. Fathizadeh, M. Khalkhali, *Scalar Curvature for the Noncommutative Two Torus* (2011), Journal of Noncommutative Geometry.

F. Fathizadeh, M. Khalkhali, *Weyl's Law and Connes' Trace Theorem for Noncommutative Two Tori* (2011), LMP.

A. Connes, H. Moscovici, *Modular Curvature for Noncommutative Two-Tori* (2011).

F. Fathizadeh, M. Khalkhali, *Scalar Curvature for Noncommutative Four-Tori*, Preprint, 2013.

F. Fathizadeh, M. Khalkhali, A. Moatadelro, *A Riemann-Roch Theorem for Noncommutative Two Torus*, Preprint, 2012.

Laplace spectrum; commutative background

- ▶ (M, g) = closed Riemannian manifold. Laplacian on forms

$$\Delta = (d + d^*)^2 : \Omega^p(M) \rightarrow \Omega^p(M),$$

has pure point spectrum:

$$0 \leq \lambda_1 \leq \lambda_2 \leq \dots \rightarrow \infty$$

- ▶ Fact: Dimension, volume, total scalar curvature, Betti numbers, and hence the Euler characteristic of M are fully determined by the spectrum of Δ (on all p -forms).

Method of proof: bring in the heat kernel

- ▶ Heat equation for functions: $\partial_t + \Delta = 0$
- ▶ $k(t, x, y) = \text{kernel of } e^{-t\Delta}$. Asymptotic expansion near $t = 0$:

$$k(t, x, x) \sim \frac{1}{(4\pi t)^{m/2}} (a_0(x, \Delta) + a_1(x, \Delta)t + a_2(x, \Delta)t^2 + \dots)$$

- ▶ $a_i(x, \Delta)$, Seeley-De Witt-Gilkey coefficients.

- ▶ Theorem: $a_i(x, \Delta)$ are universal polynomials in curvature tensor R and its covariant derivatives:

$$a_0(x, \Delta) = 1$$

$$a_1(x, \Delta) = \frac{1}{6}S(x) \quad \text{scalar curvature}$$

$$a_2(x, \Delta) = \frac{1}{360}(2|R(x)|^2 - 2|\text{Ric}(x)|^2 + 5|S(x)|^2)$$

$$a_3(x, \Delta) = \dots\dots$$

Heat trace asymptotics

Compute $\text{Trace}(e^{-t\Delta})$ in two ways:

Spectral Sum = Geometric Sum.

$$\sum e^{-t\lambda_i} = \int_M k(t, x, x) d\text{vol}_x \sim (4\pi t)^{-\frac{m}{2}} \sum_{j=0}^{\infty} a_j t^j \quad (t \rightarrow 0).$$

Hence

$$a_j = \int_M a_j(x, \Delta) d\text{vol}_x,$$

are manifestly spectral invariants:

$$a_0 = \int_M d\text{vol}_x = \text{Vol}(M), \quad \implies \text{Weyl's law}$$

$$a_1 = \frac{1}{6} \int_M S(x) d\text{vol}_x, \quad \text{total scalar curvature}$$

Tauberian theory and $a_0 = 1$, implies [Weyl's law](#):

$$N(\lambda) \sim \frac{\text{Vol}(M)}{(4\pi)^{m/2} \Gamma(1 + m/2)} \lambda^{m/2} \quad \lambda \rightarrow \infty,$$

where

$$N(\lambda) = \#\{\lambda_i \leq \lambda\}$$

is the eigenvalue counting function.

Simplest example: flat tori

- ▶ $\Gamma \subset \mathbb{R}^m$ a cocompact lattice; $M = \mathbb{R}^m/\Gamma$

$$\text{spec}(\Delta) = \{4\pi^2 \|\gamma^*\|^2; \gamma^* \in \Gamma^*\}$$

- ▶ Then:

$$K(t, x, y) = \frac{1}{(4\pi t)^{m/2}} \sum_{\gamma \in \Gamma} e^{-4\pi^2 \|x-y+\gamma\|^2/4t}$$

- ▶ Poisson summation formula \implies

$$\sum_{\gamma^* \in \Gamma^*} e^{-4\pi^2 \|\gamma^*\|^2 t} = \frac{\text{Vol}(M)}{(4\pi t)^{m/2}} \sum_{\gamma \in \Gamma} e^{-4\pi^2 \|\gamma\|^2/4t}$$

- ▶ And from this we obtain the asymptotic expansion of the heat trace near $t = 0$

$$\text{Tr} e^{-t\Delta} \sim \frac{\text{Vol}(M)}{(4\pi t)^{m/2}} \quad (t \rightarrow 0)$$

Application 1: heat equation proof of the Atiyah-Singer index theorem

- ▶ Dirac operator

$$D : C^\infty(S_+) \rightarrow C^\infty(S_-)$$

McKean-Singer formula:

$$\text{Index}(D) = \text{Tr}(e^{-tD^*D}) - \text{Tr}(e^{-tDD^*}), \quad \forall t > 0$$

Heat trace asymptotics \implies

$$\text{Index}(D) = \int_M a_n(x) dx,$$

where $a_n(x) = a_n^+(x) - a_n^-(x)$, $m = 2n$, can be explicitly computed and recovers the A-S integrand (The simplest proof is due to Getzler).

Application 2: meromorphic extension of spectral zeta functions

$$\zeta_{\Delta}(s) := \sum_{\lambda_j \neq 0} \lambda_j^{-s}, \quad \operatorname{Re}(s) > \frac{m}{2}$$

Mellin transform + asymptotic expansion:

$$\lambda^{-s} = \frac{1}{\Gamma(s)} \int_0^{\infty} e^{-t} t^{s-1} dt \quad \operatorname{Re}(s) > 0$$

$$\begin{aligned} \zeta_{\Delta}(s) &= \frac{1}{\Gamma(s)} \int_0^{\infty} (\operatorname{Trace}(e^{-t\Delta}) - \dim \operatorname{Ker} \Delta) t^{s-1} dt \\ &= \frac{1}{\Gamma(s)} \left\{ \int_0^c \dots + \int_c^{\infty} \dots \right\} \end{aligned}$$

The second term defines an entire function, while the first term has a meromorphic extension to \mathbb{C} with **simple poles** within the set

$$\frac{m}{2} - j, \quad j = 0, 1, \dots$$

Also: 0 is always a regular point.

Simplest example: For $M = S^1$ with round metric, we have

$$\zeta_{\Delta}(s) = 2\zeta(2s) \quad \text{Riemann zeta function}$$

Scalar curvature

The spectral invariants a_j in the heat asymptotic expansion

$$\text{Trace}(e^{-t\Delta}) \sim (4\pi t)^{-\frac{m}{2}} \sum_{j=0}^{\infty} a_j t^j \quad (t \rightarrow 0)$$

are related to residues of spectral zeta function by

$$\text{Res}_{s=\alpha} \zeta_{\Delta}(s) = (4\pi)^{-\frac{m}{2}} \frac{a_{\frac{m}{2}-\alpha}}{\Gamma(\alpha)}, \quad \alpha = \frac{m}{2} - j > 0$$

Focusing on subleading pole $s = \frac{m}{2} - 1$ and using $a_1 = \frac{1}{6} \int_M S(x) d\text{vol}_x$, we obtain a formula for scalar curvature density as follows:

Let $\zeta_f(s) := \text{Tr}(f\Delta^{-s})$, $f \in C^\infty(M)$.

$$\text{Res } \zeta_f(s)|_{s=\frac{m}{2}-1} = \frac{(4\pi)^{-m/2}}{\Gamma(m/2-1)} \int_M fS(x) d\text{vol}_x, \quad m \geq 3$$

$$\zeta_f(s)|_{s=0} = \frac{1}{4\pi} \int_M fS(x) d\text{vol}_x - \text{Tr}(fP) \quad m = 2$$

$\log \det(\Delta) = -\zeta'(0)$, Ray-Singer regularized determinant

Noncommutative Geometry: Spectral Triples $(\mathcal{A}, \mathcal{H}, D)$

- ▶ \mathcal{A} = involutive unital algebra, \mathcal{H} = Hilbert space,

$$\pi : \mathcal{A} \rightarrow \mathcal{L}(\mathcal{H}), \quad D : \mathcal{H} \rightarrow \mathcal{H}$$

D has compact resolvent and all commutators $[D, \pi(a)]$ are bounded.

- ▶ An asymptotic expansion holds

$$\text{Trace}(e^{-tD^2}) \sim \sum a_\alpha t^\alpha \quad (t \rightarrow 0)$$

- ▶ Let $\Delta = D^2$. Spectral zeta function

$$\zeta_D(s) = \text{Tr}(|D|^{-s}) = \text{Tr}(\Delta^{-s/2}), \quad \text{Re}(s) \gg 0.$$

Noncommutative 4-torus \mathbb{T}_θ^4

$C(\mathbb{T}_\theta^4)$ is the universal C^* -algebra generated by 4 unitaries

$$U_1, U_2, U_3, U_4,$$

satisfying

$$U_k U_\ell = e^{2\pi i \theta_{k\ell}} U_\ell U_k,$$

for a skew symmetric matrix

$$\theta = (\theta_{k\ell}) \in M_4(\mathbb{R}).$$

Action of $\mathbb{T}^4 = (\mathbb{R}/2\pi\mathbb{Z})^4$ on $C(\mathbb{T}_\theta^4)$

$$\mathbb{R}^4 \ni \mathbf{s} \mapsto \alpha_{\mathbf{s}} \in \text{Aut}\left(C(\mathbb{T}_\theta^4)\right),$$

$$\alpha_{\mathbf{s}}(U^{\mathbf{m}}) := e^{i\mathbf{s} \cdot \mathbf{m}} U^{\mathbf{m}}, \quad U^{\mathbf{m}} := U_1^{m_1} U_2^{m_2} U_3^{m_3} U_4^{m_4}, \quad m_j \in \mathbb{Z}.$$

$$\delta_j = \frac{\partial}{\partial s_j} \Big|_{s=0} \alpha_{\mathbf{s}} : C^\infty(\mathbb{T}_\theta^4) \rightarrow C^\infty(\mathbb{T}_\theta^4),$$

$$\begin{aligned} \delta_j(U_k) &:= U_k && \text{if } k = j, \\ &:= 0 && \text{if } k \neq j. \end{aligned}$$

Complex Structure on \mathbb{T}_θ^4

$$\partial = \partial_1 \oplus \partial_2, \quad \bar{\partial} = \bar{\partial}_1 \oplus \bar{\partial}_2,$$

$$\partial_1 = \frac{1}{2}(\delta_1 - i\delta_3), \quad \partial_2 = \frac{1}{2}(\delta_2 - i\delta_4),$$

$$\bar{\partial}_1 = \frac{1}{2}(\delta_1 + i\delta_3), \quad \bar{\partial}_2 = \frac{1}{2}(\delta_2 + i\delta_4).$$

Volume form on \mathbb{T}_θ^4

$$\varphi_0 : C(\mathbb{T}_\theta^4) \rightarrow \mathbb{C},$$

$$\varphi_0(1) := 1,$$

$$\varphi_0(U_1^{m_1} U_2^{m_2} U_3^{m_3} U_4^{m_4}) := 0, \quad (m_1, m_2, m_3, m_4) \neq (0, 0, 0, 0).$$

$$\varphi_0(ab) = \varphi_0(ba), \quad a, b \in C(\mathbb{T}_\theta^4).$$

$$\varphi_0(a^* a) > 0, \quad a \neq 0.$$

Conformal perturbation (Connes-Tretkoff)

Let $h = h^* \in C^\infty(\mathbb{T}_\theta^4)$ and replace the trace φ_0 by

$$\varphi : C(\mathbb{T}_\theta^4) \rightarrow \mathbb{C},$$

$$\varphi(a) := \varphi_0(a e^{-2h}), \quad a \in C(\mathbb{T}_\theta^4).$$

φ is a KMS state with the modular group

$$\sigma_t(a) = e^{2ith} a e^{-2ith}, \quad a \in C(\mathbb{T}_\theta^4),$$

and the modular automorphism

$$\Delta(a) := \sigma_i(a) = e^{-2h} a e^{2h}, \quad a \in C(\mathbb{T}_\theta^4).$$

$$\varphi(ab) = \varphi(b \Delta(a)), \quad a, b \in C(\mathbb{T}_\theta^4).$$

Perturbed Laplacian on \mathbb{T}_θ^4

$$d = \partial \oplus \bar{\partial} : \mathcal{H}_\varphi \rightarrow \mathcal{H}_\varphi^{(1,0)} \oplus \mathcal{H}_\varphi^{(0,1)},$$

$$\Delta_\varphi := d^* d.$$

Remark. If $h = 0$ then $\varphi = \varphi_0$ and

$$\Delta_{\varphi_0} = \delta_1^2 + \delta_2^2 + \delta_3^2 + \delta_4^2 = \partial^* \partial$$

(the underlying manifold is Kähler).

Scalar Curvature for \mathbb{T}_θ^4

It is the unique element $R \in C^\infty(\mathbb{T}_\theta^4)$ such that

$$\operatorname{Res}_{s=1} \zeta_a(s) = \varphi_0(a R), \quad a \in C^\infty(\mathbb{T}_\theta^4),$$

$$\zeta_a(s) := \operatorname{Trace}(a \Delta_\varphi^{-s}), \quad \Re(s) \gg 0.$$

Lemma. Up to an anti-unitary equivalence Δ_φ is given by

$$e^h \bar{\partial}_1 e^{-h} \partial_1 e^h + e^h \partial_1 e^{-h} \bar{\partial}_1 e^h + e^h \bar{\partial}_2 e^{-h} \partial_2 e^h + e^h \partial_2 e^{-h} \bar{\partial}_2 e^h,$$

where $\partial_1, \bar{\partial}_2$ are analogues of the Dolbeault operators.

Mellin Transform and Asymptotic Expansions

$$\text{Trace}(a e^{-t\Delta_\varphi}) \sim_{t \rightarrow 0^+} t^{-2} \sum_{n=0}^{\infty} B_n(a, \Delta_\varphi) t^{n/2}.$$

Approximate $e^{-t\Delta_\varphi}$ by pseudodifferential operators:

$$e^{-t\Delta_\varphi} = \frac{1}{2\pi i} \int_C e^{-t\lambda} (\Delta_\varphi - \lambda)^{-1} d\lambda,$$

$$B_\lambda (\Delta_\varphi - \lambda) \approx 1,$$

$$\sigma(B_\lambda) = b_0 + b_1 + b_2 + \dots .$$

Connes' Pseudodifferential Calculus (1980)

- Symbols: $\rho : \mathbb{R}^4 \rightarrow C^\infty(\mathbb{T}_\theta^4)$,

$$P_\rho(a) = (2\pi)^{-4} \int \int e^{-is \cdot \xi} \rho(\xi) \alpha_s(a) ds d\xi, \quad a \in C^\infty(\mathbb{T}_\theta^4).$$

- Differential operators:

$$\rho(\xi) = \sum a_\ell \xi^\ell, \quad a_\ell \in C^\infty(\mathbb{T}_\theta^4) \Rightarrow P_\rho = \sum a_\ell \delta^\ell.$$

- Ψ DO's on \mathbb{T}_θ^4 form an algebra:

$$\sigma(PQ) \sim \sum_{\ell \in \mathbb{Z}_{\geq 0}^4} \frac{1}{\ell!} \partial_\xi^\ell \rho(\xi) \delta^\ell(\rho'(\xi)).$$

Connes' Rearrangement Lemma

For any $m = (m_0, m_1, \dots, m_\ell) \in \mathbb{Z}_{>0}^{\ell+1}$, $\rho_1, \dots, \rho_\ell \in C^\infty(\mathbb{T}_\theta^4)$

$$\begin{aligned} & \int_0^\infty \frac{u^{|m|-2}}{(e^h u + 1)^{m_0}} \prod_1^\ell \rho_j (e^h u + 1)^{-m_j} du \\ &= e^{-(|m|-1)h} F_m(\Delta, \dots, \Delta) \left(\prod_1^\ell \rho_j \right), \end{aligned}$$

where

$$F_m(u_1, \dots, u_\ell) = \int_0^\infty \frac{x^{|m|-2}}{(x+1)^{m_0}} \prod_1^\ell \left(x \prod_1^j u_k + 1 \right)^{-m_j} dx.$$

Final Formula for the Scalar Curvature of \mathbb{T}_θ^4

Theorem.

$$R = e^{-h} k(\nabla) \left(\sum_{i=1}^4 \delta_i^2(h) \right) + e^{-h} H(\nabla, \nabla) \left(\sum_{i=1}^4 \delta_i(h)^2 \right),$$

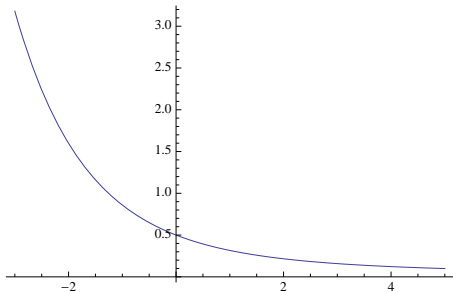
where

$$\nabla(a) = [-h, a], \quad a \in C(\mathbb{T}_\theta^4),$$

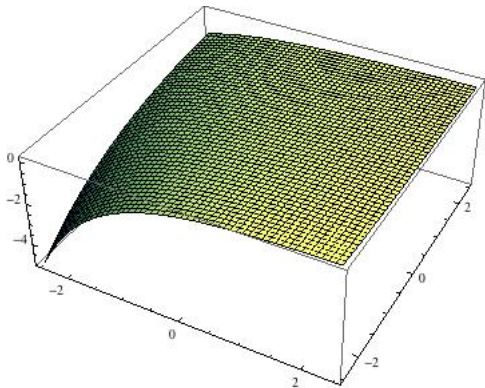
$$k(s) = \frac{1 - e^{-s}}{2s},$$

$$H(s, t) = -\frac{e^{-s-t} ((-e^s - 3)s(e^t - 1) + (e^s - 1)(3e^t + 1)t)}{4st(s+t)}.$$

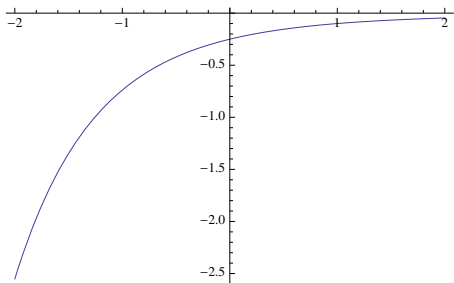
$$k(s) = \frac{1}{2} - \frac{s}{4} + \frac{s^2}{12} - \frac{s^3}{48} + \frac{s^4}{240} - \frac{s^5}{1440} + O(s^6).$$



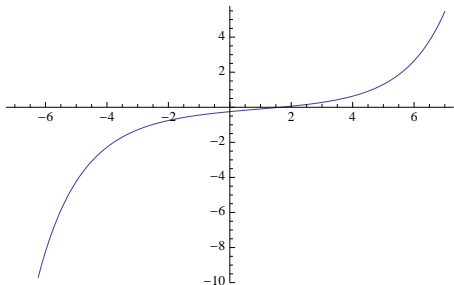
$$H(s, t) = \left(-\frac{1}{4} + \frac{t}{24} + O(t^3) \right) + s \left(\frac{5}{24} - \frac{t}{16} + \frac{t^2}{80} + O(t^3) \right) \\ + s^2 \left(-\frac{1}{12} + \frac{7t}{240} - \frac{t^2}{144} + O(t^3) \right) + O(s^3).$$



$$\begin{aligned}
 H(s, s) &= -\frac{e^{-2s}(e^s - 1)^2}{4s^2} \\
 &= -\frac{1}{4} + \frac{s}{4} - \frac{7s^2}{48} + \frac{s^3}{16} - \frac{31s^4}{1440} + \frac{s^5}{160} + O(s^6).
 \end{aligned}$$



$$\begin{aligned}
 G(s) &:= H(s, -s) = \frac{-4s - 3e^{-s} + e^s + 2}{4s^2} \\
 &= -\frac{1}{4} + \frac{s}{6} - \frac{s^2}{48} + \frac{s^3}{120} - \frac{s^4}{1440} + \frac{s^5}{5040} + O(s^6).
 \end{aligned}$$



Einstein-Hilbert Action for \mathbb{T}_θ^4

Theorem. We have the local expression (up to a factor of π^2)

$$\begin{aligned}\varphi_0(R) &= \frac{1}{2} \sum_{i=1}^4 \varphi_0\left(e^{-h} \delta_i^2(h)\right) \\ &\quad + \sum_{i=1}^4 \varphi_0\left(G(\nabla)(e^{-h} \delta_i(h)) \delta_i(h)\right).\end{aligned}$$

Extremum of the Einstein-Hilbert Action

Theorem. For any Weyl factor $e^{-h} \in C^\infty(\mathbb{T}_\theta^4)$

$$\varphi_0(R) \leq 0,$$

and the equality happens if and only if h is a constant.

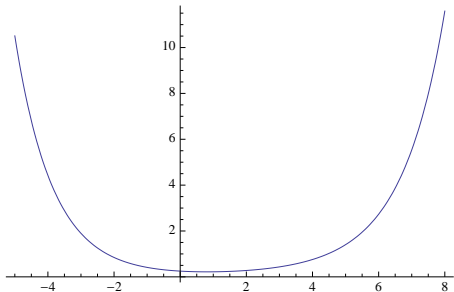
Proof.

$$\varphi_0(R) = \sum_{i=1}^4 \varphi_0(e^{-h} T(\nabla)(\delta_i(h)) \delta_i(h)),$$

where

$$T(s) = \frac{1}{2} \frac{e^{-s} - 1}{-s} + G(s) = \frac{-2s + e^s - e^{-s}(2s + 3) + 2}{4s^2}.$$

$$T(s) = \frac{1}{4} - \frac{s}{12} + \frac{s^2}{16} - \frac{s^3}{80} + \frac{s^4}{288} - \frac{s^5}{2016} + O(s^6).$$



Weyl's Law for \mathbb{T}_θ^4

Theorem. For the eigenvalue counting function

$$N(\lambda) = \#\{\lambda_j \leq \lambda\}$$

of the Laplacian Δ_φ on \mathbb{T}_θ^4 , we have

$$N(\lambda) \sim \frac{\pi^2 \varphi_0(e^{-2h})}{2} \lambda^2 \quad (\lambda \rightarrow \infty).$$

Corollary.

$$\lambda_j \sim \frac{\sqrt{2}}{\pi \varphi_0(e^{-2h})^{1/2}} j^{1/2} \quad (j \rightarrow \infty),$$

$$\mathrm{Tr}_\omega \left((1 + \Delta_\varphi)^{-2} \right) = \frac{\pi^2}{2} \varphi_0(e^{-2h}).$$

A Noncommutative Residue

Classical symbols: $\rho : \mathbb{R}^4 \rightarrow A_\theta^\infty$

$$\rho(\xi) \sim \sum_{i=0}^{\infty} \rho_{m-i}(\xi) \quad (\xi \rightarrow \infty),$$

$$\rho_{m-i}(t\xi) = t^{m-i} \rho_{m-i}(\xi), \quad t > 0, \quad \xi \in \mathbb{R}^4.$$

$$\text{Res}(P_\rho) := \int_{\mathbb{S}^3} \varphi_0(\rho_{-4}(\xi)) d\xi.$$

Theorem. Res is the unique trace on the algebra of classical pseudodifferential operators on \mathbb{T}_θ^4 .

Connes' Trace Theorem for \mathbb{T}_θ^4

Theorem. For any classical symbol ρ of order -4 on \mathbb{T}_θ^4 , we have

$$P_\rho \in \mathcal{L}^{1,\infty}(\mathcal{H}_0),$$

and

$$\mathrm{Tr}_\omega(P_\rho) = \frac{1}{4} \mathrm{Res}(P_\rho).$$

The geometry in noncommutative geometry

- ▶ Geometry starts with **metric** and **curvature**. While there are a good number of ‘soft’ topological tools in NCG, like cyclic cohomology, K and KK-theory, and index theory, a truly noncommutative theory of curvature is still illusive. The situation is better with **scalar curvature**, but computations are quite tough at the moment.
- ▶ Metric aspects of NCG are informed by **Spectral Geometry**. Spectral invariants are the only means by which we can formulate metric ideas of NCG.