

Spectral Zeta Functions and Scalar Curvature for Noncommutative Tori

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- ▶ A. Connes and P. Tretkoff, *The Gauss-Bonnet Theorem for the noncommutative two torus* , 1992, and Sept. 2009.
- ▶ A. Connes and H. Moscovici, *Modular curvature for noncommutative two-tori*, Oct. 2011.
- ▶ F. Fathizadeh and M. Khalkhali, *The Gauss-Bonnet Theorem for noncommutative two tori with a general conformal structure*, May 2010.
- ▶ F. Fathizadeh and M. Khalkhali, *Scalar Curvature for the Noncommutative Two Torus*, Oct. 2011.

From heat trace and zeta functions to spectral invariants

(M, g) = closed Riemannian manifold. [Laplacian on functions](#)

$$\Delta = d^*d : C^\infty(M) \rightarrow C^\infty(M)$$

It is positive with pure point spectrum. Let $k(t, x, y)$ = kernel of $e^{-t\Delta}$:

$$k(t, x, y) \sim \frac{1}{(4\pi t)^{m/2}} e^{-d(x,y)^2/4t}, \quad (t \rightarrow 0)$$

Restrict to the diagonal: as $t \rightarrow 0$,

$$k(t, x, x) \sim \frac{1}{(4\pi t)^{m/2}} (a_0(x, \Delta) + a_1(x, \Delta)t + a_2(x, \Delta)t^2 + \dots)$$

- ▶ Functions $a_i(x, \Delta)$: expressed by universal polynomials in curvature tensor R and its covariant derivatives:

$$a_0(x, \Delta) = 1$$

$$a_1(x, \Delta) = \frac{1}{6}S(x) \quad \text{scalar curvature}$$

$$a_2(x, \Delta) = \frac{1}{360}(2|R(x)|^2 - 2|\text{Ric}(x)|^2 + 5|S(x)|^2)$$

$$a_3(x, \Delta) = \dots\dots\dots$$

Short time asymptotics of the heat trace

$$\begin{aligned}\text{Trace}(e^{-t\Delta}) &= \sum e^{-t\lambda_i} = \int_M k(t, x, x) d\text{vol}_x \\ &\sim (4\pi t)^{-\frac{m}{2}} \sum_{j=0}^{\infty} a_j t^j \quad (t \rightarrow 0)\end{aligned}$$

So

$$a_j = \int_M a_j(x, \Delta) d\text{vol}_x,$$

are manifestly spectral invariants.

$$\begin{aligned}a_0 &= \int_M d\text{vol}_x = \text{Vol}(M), && \implies \text{Weyl's law} \\ a_1 &= \frac{1}{6} \int_M S(x) d\text{vol}_x, && \text{total scalar curvature}\end{aligned}$$

Spectral zeta functions

$$\zeta_{\Delta}(s) := \sum_{\lambda_j \neq 0} \lambda_j^{-s}, \quad \operatorname{Re}(s) > \frac{m}{2}$$

Mellin transform + asymptotic expansion \implies

$$\zeta_{\Delta}(s) = \frac{1}{\Gamma(s)} \int_0^{\infty} (\operatorname{Trace}(e^{-t\Delta}) - \operatorname{Dim Ker} \Delta) t^{s-1} dt$$

has meromorphic extension to \mathbb{C} with **simple poles** within the set

$$\frac{m}{2} - j, \quad j = 0, 1, \dots$$

Also: 0 is always a regular point.

Let $\zeta_f(s) := \text{Tr}(f\Delta^{-s})$, $f \in C^\infty(M)$.

$$\text{Res } \zeta_f(s)|_{s=\frac{m}{2}-1} = \frac{(4\pi)^{-m/2}}{\Gamma(m/2-1)} \int_M fS(x) d\text{vol}_x, \quad m \geq 3$$

$$\zeta_f(s)|_{s=0} = \frac{1}{4\pi} \int_M fS(x) d\text{vol}_x - \text{Tr}(fP) \quad m = 2$$

$$\log \det(\Delta) = -\zeta'(0), \quad \text{regularized determinant}$$

Spectral Triples: $(\mathcal{A}, \mathcal{H}, D)$

- ▶ \mathcal{A} = involutive unital algebra, \mathcal{H} = Hilbert space,

$$\pi : \mathcal{A} \rightarrow \mathcal{L}(\mathcal{H}), \quad D : \mathcal{H} \rightarrow \mathcal{H}$$

D has compact resolvent and all commutators $[D, \pi(a)]$ are bounded.

- ▶ Assume: an asymptotic expansion of the form

$$\text{Trace}(e^{-tD^2}) \sim \sum a_\alpha t^\alpha \quad (t \rightarrow 0)$$

holds.

- ▶ Let $\Delta = D^2$. Spectral zeta function

$$\zeta_D(s) = \text{Tr}(|D|^{-s}) = \text{Tr}(\Delta^{-s/2}), \quad \text{Re}(s) \gg 0.$$

- ▶ Using the Mellin transform and the asymptotic expansion, easy to show that: ζ_D has a meromorphic extension to all of \mathbb{C} and non-zero terms, a_α , $\alpha < 0$, give a pole of ζ_D at -2α with

$$\operatorname{Res}_{s=-2\alpha} \zeta_D(s) = \frac{2a_\alpha}{\Gamma(-\alpha)}.$$

Also, $\zeta_D(s)$ is holomorphic at $s = 0$ and

$$\zeta_D(0) + \dim \ker D = a_0$$

Curved noncommutative tori A_θ , $\theta \in \mathbb{R}$

A_θ = universal C^* -algebra generated by unitaries U and V

$$VU = e^{2\pi i\theta} UV.$$

$$A_\theta^\infty = \left\{ \sum_{m,n \in \mathbb{Z}} a_{m,n} U^m V^n : (a_{m,n}) \text{ Schwartz class} \right\}.$$

► Differential operators on A_θ

$$\delta_1, \delta_2 : A_\theta^\infty \rightarrow A_\theta^\infty,$$

Infinitesimal generators of the action

$$\alpha_s(U^m V^n) = e^{is \cdot (m,n)} U^m V^n \quad s \in \mathbb{R}^2.$$

Analogues of $\frac{1}{i} \frac{\partial}{\partial x}$, $\frac{1}{i} \frac{\partial}{\partial y}$ on 2-torus.

► Canonical trace $\mathfrak{t} : A_\theta \rightarrow \mathbb{C}$ on smooth elements:

$$\mathfrak{t}\left(\sum_{m,n \in \mathbb{Z}} a_{m,n} U^m V^n\right) = a_{0,0}.$$

Complex structures on A_θ

► $\tau = \tau_1 + i\tau_2, \quad \tau_2 = \Im(\tau) > 0,$

$$\partial := \delta_1 + \tau\delta_2, \quad \partial^* := \delta_1 + \bar{\tau}\delta_2.$$

- Hilbert space of $(1, 0)$ -forms:

$\mathcal{H}^{(1,0)} :=$ completion of finite sums $\sum a\partial b, a, b \in A_\theta^\infty$, w.r.t.

$$\langle a\partial b, a'\partial b' \rangle := \mathfrak{t}((a'\partial b')^* a\partial b).$$

- Flat Dolbeault Laplacian: $\partial^*\partial = \delta_1^2 + 2\tau_1\delta_1\delta_2 + |\tau|^2\delta_2^2.$

Conformal perturbation of a metric

- ▶ Fix $h = h^* \in A_\theta^\infty$. Replace the volume form \mathfrak{t} by $\varphi : A_\theta \rightarrow \mathbb{C}$,

$$\varphi(a) := \mathfrak{t}(ae^{-h}), \quad a \in A_\theta.$$

- ▶ It is a KMS state with **modular group**

$$\sigma_t(x) = e^{ith} x e^{-ith},$$

and **modular automorphism** (Tomita-Takesaki theory)

$$\sigma_i(x) = \Delta(x) = e^{-h} x e^h.$$

$$\varphi(ab) = \varphi(b\Delta(a)), \quad \forall a, b \in A_\theta.$$

- ▶ Warning: Δ and Δ are very different operators!

Connes-Tretkoff spectral triple

- ▶ Hilbert space $\mathcal{H}_\varphi :=$ completion of A_θ w.r.t. $\langle \cdot, \cdot \rangle_\varphi$,

$$\langle a, b \rangle_\varphi := \varphi(b^* a), \quad a, b \in A_\theta$$

(GNS construction).

- ▶ View

$$\partial_\varphi = \partial = \delta_1 + \tau\delta_2 : \mathcal{H}_\varphi \rightarrow \mathcal{H}^{(1,0)}.$$

$$\mathcal{H} = \mathcal{H}_\varphi \oplus \mathcal{H}^{(1,0)},$$

$$D = \begin{pmatrix} 0 & \partial_\varphi^* \\ \partial_\varphi & 0 \end{pmatrix} : \mathcal{H} \rightarrow \mathcal{H}.$$

Full perturbed Laplacian:

$$\Delta := D^2 = \begin{pmatrix} \partial_\varphi^* \partial_\varphi & 0 \\ 0 & \partial_\varphi \partial_\varphi^* \end{pmatrix} : \mathcal{H} \rightarrow \mathcal{H}.$$

Lemma: $\partial_\varphi^* \partial_\varphi : \mathcal{H}_\varphi \rightarrow \mathcal{H}_\varphi$, and $\partial_\varphi \partial_\varphi^* : \mathcal{H}^{(1,0)} \rightarrow \mathcal{H}^{(1,0)}$ are anti-unitarily equivalent to

$$\begin{aligned} k \partial^* \partial k &: \mathcal{H}_0 \rightarrow \mathcal{H}_0, \\ \partial^* k^2 \partial &: \mathcal{H}^{(1,0)} \rightarrow \mathcal{H}^{(1,0)}, \end{aligned}$$

where $k = e^{h/2}$.

The Tomita anti-unitary map J is used.

Scalar curvature for A_θ

- ▶ The scalar curvature of the curved nc torus $(\mathbb{T}_\theta^2, \tau, k)$ is the unique element $R \in A_\theta^\infty$ satisfying

$$\text{Trace}(a\Delta^{-s})|_{s=0} + \text{Trace}(aP) = \text{t}(aR), \quad \forall a \in A_\theta^\infty,$$

where P is the projection onto the kernel of Δ .

- ▶ The **chiral scalar curvature** R^γ is the unique element $R^\gamma \in A_\theta^\infty$ which satisfies the equation

$$\text{Trace}(\gamma a \Delta^{-s})|_{s=0} + \text{t}(ae^{-h})/\text{t}(e^{-h}) - \text{t}(a) = \text{t}(aR^\gamma), \quad \forall a \in A_\theta^\infty.$$

- ▶ We find a formula for R using **Connes' pseudodifferential calculus**.

Connes' pseudodifferential calculus

- ▶ Symbols: $\rho : \mathbb{R}^2 \rightarrow A_\theta^\infty$.
- ▶ Ψ DO's: $P_\rho : A_\theta^\infty \rightarrow A_\theta^\infty$,

$$P_\rho(a) = (2\pi)^{-2} \int_{\mathbb{R}^2} \int_{\mathbb{R}^2} e^{-is \cdot \xi} \rho(\xi) \alpha_s(a) ds d\xi.$$

- ▶ For example:

$$\rho(\xi_1, \xi_2) = \sum a_{ij} \xi_1^i \xi_2^j, \quad a_{ij} \in A_\theta^\infty \quad \Rightarrow \quad P_\rho = \sum a_{ij} \delta_1^i \delta_2^j.$$

- ▶ Multiplication of symbol.

$$\sigma(PQ) \sim \sum_{\ell_1, \ell_2 \geq 0} \frac{1}{\ell_1! \ell_2!} \partial_1^{\ell_1} \partial_2^{\ell_2} (\rho(\xi)) \delta_1^{\ell_1} \delta_2^{\ell_2} (\rho'(\xi)).$$

Local expression for the scalar curvature

- ▶ Cauchy integral formula:

$$e^{-t\Delta} = \frac{1}{2\pi i} \int_C e^{-t\lambda} (\Delta - \lambda)^{-1} d\lambda.$$

- ▶ $B_\lambda \approx (\Delta - \lambda)^{-1}$:

$$\sigma(B_\lambda) \sim b_0(\xi, \lambda) + b_1(\xi, \lambda) + b_2(\xi, \lambda) + \dots,$$

each $b_j(\xi, \lambda)$ is a symbol of order $-2 - j$, and

$$\sigma(B_\lambda(\Delta - \lambda)) \sim 1.$$

(Note: λ is considered of order 2.)

Proposition: The scalar curvature of the spectral triple attached to (A_θ, τ, k) is equal to

$$\frac{1}{2\pi i} \int_{\mathbb{R}^2} \int_C e^{-\lambda} b_2(\xi, \lambda) d\lambda d\xi,$$

where b_2 is defined as above.

The computations for $k\partial^*\partial k$

- ▶ The symbol of $k\partial^*\partial k$ is equal to

$$a_2(\xi) + a_1(\xi) + a_0(\xi)$$

where

$$a_2(\xi) = \xi_1^2 k^2 + |\tau|^2 \xi_2^2 k^2 + 2\tau_1 \xi_1 \xi_2 k^2,$$

$$a_1(\xi) = 2\xi_1 k \delta_1(k) + 2|\tau|^2 \xi_2 k \delta_2(k) + 2\tau_1 \xi_1 k \delta_2(k) + 2\tau_1 \xi_2 k \delta_1(k),$$

$$a_0(\xi) = k \delta_1^2(k) + |\tau|^2 k \delta_2^2(k) + 2\tau_1 k \delta_1 \delta_2(k).$$

- ▶ The equation

$$(b_0 + b_1 + b_2 + \dots)((a_2 + 1) + a_1 + a_0) \sim 1,$$

has a solution with each b_j a symbol of order $-2 - j$.

$$b_0 = (a_2 + 1)^{-1} = (\xi_1^2 k^2 + |\tau|^2 \xi_2^2 k^2 + 2\tau_1 \xi_1 \xi_2 k^2 + 1)^{-1},$$

$$b_1 = -(b_0 a_1 b_0 + \partial_1(b_0) \delta_1(a_2) b_0 + \partial_2(b_0) \delta_2(a_2) b_0),$$

$$b_2 = -(b_0 a_0 b_0 + b_1 a_1 b_0 + \partial_1(b_0) \delta_1(a_1) b_0 + \partial_2(b_0) \delta_2(a_1) b_0 + \\ \partial_1(b_1) \delta_1(a_2) b_0 + \partial_2(b_1) \delta_2(a_2) b_0 + (1/2) \partial_{11}(b_0) \delta_1^2(a_2) b_0 + \\ (1/2) \partial_{22}(b_0) \delta_2^2(a_2) b_0 + \partial_{12}(b_0) \delta_{12}(a_2) b_0)$$

$$= 5\xi_1^2 b_0^2 k^3 \delta_1^2(k) b_0 + 2\xi_1^2 b_0 k \delta_1(k) b_0 \delta_1(k) b_0 k \\ + \text{about 800 terms.}$$

To integrate b_2 over the ξ -plane, pass to the coordinates

$$\xi_1 = r \cos \theta - r \frac{\tau_1}{\tau_2} \sin \theta, \quad \xi_2 = \frac{r}{\tau_2} \sin \theta.$$

After the integration with respect to θ , up to a factor of $\frac{r}{\tau_2}$ which is the Jacobian of the change of variables, one gets

$$4|\tau|^2 \pi r^6 b_0^2 k^2 \delta_2(k) b_0^2 k^3 \delta_2(k) b_0 k + 4\tau_1 \pi r^6 b_0^2 k^2 \delta_2(k) b_0^2 k^3 \delta_1(k) b_0 k \\ + \dots \quad (\text{more than 80 terms})$$

Terms with two b_0^i factors

These are the following terms

$$-4\pi r^4 b_0^3 k^4 \delta_1^2(k) b_0 k - 4|\tau|^2 \pi r^4 b_0^3 k^4 \delta_2^2(k) b_0 k + \dots (23 \text{ terms})$$

where

$$b_0 = (r^2 k^2 + 1)^{-1}.$$

Rearrangement lemma (Connes)

The computation of $\int_0^\infty \bullet r dr$ of these terms is achieved by:

For all $\rho_j \in A_\theta^\infty$ and $m_j > 0$ one has

$$\begin{aligned} & \int_0^\infty (k^2 u + 1)^{-m_0} \rho_1(k^2 u + 1)^{-m_1} \cdots \rho_\ell(k^2 u + 1)^{-m_\ell} u^{\sum m_j - 2} du \\ &= k^{-2(\sum m_j - 1)} F_{m_0, m_1, \dots, m_\ell}(\Delta_{(1)}, \Delta_{(2)}, \dots, \Delta_{(\ell)})(\rho_1 \rho_2 \cdots \rho_\ell), \end{aligned}$$

where

$$F_{m_0, m_1, \dots, m_\ell}(u_1, u_2, \dots, u_\ell) = \int_0^\infty (u + 1)^{-m} \prod_1^\ell (u \prod_1^j u_h + 1)^{-m_j} u^{\sum m_j - 2}$$

and $\Delta_{(i)}$ signifies that Δ acts on the i -th factor.

Up to an overall factor of π , $\int_0^\infty \bullet r dr$ of the terms with two positive powers of b_0 is equal to

$$\begin{aligned}
 & f_1(\Delta)(k^{-1}\delta_1^2(k)) + f_2(\Delta)(k^{-2}\delta_1(k)^2) \\
 + & |\tau|^2 f_1(\Delta)(k^{-1}\delta_2^2(k)) + |\tau|^2 f_2(\Delta)(k^{-2}\delta_2(k)^2) \\
 + & \tau_1 f_1(\Delta)(k^{-1}\delta_1\delta_2(k)) + \tau_1 f_2(\Delta)(k^{-2}\delta_1(k)\delta_2(k)) \\
 + & \tau_1 f_1(\Delta)(k^{-1}\delta_2\delta_1(k)) + \tau_1 f_2(\Delta)(k^{-2}\delta_2(k)\delta_1(k)),
 \end{aligned}$$

where

$$\begin{aligned}
 f_1(u) & := -2\mathcal{L}_2(u)u^{1/2} - 2\mathcal{L}_2(u) + \mathcal{L}_1(u)u^{1/2} + 3\mathcal{L}_1(u) - \mathcal{L}_0(u) \\
 & = -\frac{u^{1/2}(2 - 2u + (1 + u)\log u)}{(-1 + u^{1/2})^3(1 + u^{1/2})^2},
 \end{aligned}$$

and

$$f_2(u) := -4\mathcal{L}_2(u) + 4\mathcal{L}_1(u) = 2\frac{-1 + u^2 - 2u\log u}{(-1 + u)^3}.$$

Terms with three b_0^i

These terms are the following:

$$4|\tau|^2 \pi r^6 b_0^2 k^2 \delta_2(k) b_0^2 k^3 \delta_2(k) b_0 k + 4\tau_1 \pi r^6 b_0^2 k^2 \delta_2(k) b_0^2 k^3 \delta_1(k) b_0 k \\ + \dots \quad (62 \text{ terms})$$

For computing $\int_0^\infty \bullet r dr$ of these terms we again use the *Rearrangement Lemma*.

After the integrations, up to an overall factor of π , we find the following expression

$$\begin{aligned}
 & f_1(\Delta)(k^{-1}\delta_1^2(k)) + f_2(\Delta)(k^{-2}\delta_1(k)^2) \\
 + & F(\Delta_{(1)}, \Delta_{(2)})((\delta_1(k)k^{-1})(k^{-1}\delta_1(k))) \\
 + & |\tau|^2 f_1(\Delta)(k^{-1}\delta_2^2(k)) + |\tau|^2 f_2(\Delta)(k^{-2}\delta_2(k)^2) \\
 + & |\tau|^2 F(\Delta_{(1)}, \Delta_{(2)})((\delta_2(k)k^{-1})(k^{-1}\delta_2(k))) \\
 + & \tau_1 f_1(\Delta)(k^{-1}\delta_1\delta_2(k)) + \tau_1 f_2(\Delta)(k^{-2}\delta_1(k)\delta_2(k)) \\
 + & \tau_1 F(\Delta_{(1)}, \Delta_{(2)})((\delta_1(k)k^{-1})(k^{-1}\delta_2(k))) \\
 + & \tau_1 f_1(\Delta)(k^{-1}\delta_2\delta_1(k)) + \tau_1 f_2(\Delta)(k^{-2}\delta_2(k)\delta_1(k)) \\
 + & \tau_1 F(\Delta_{(1)}, \Delta_{(2)})((\delta_2(k)k^{-1})(k^{-1}\delta_1(k))),
 \end{aligned}$$

where we have

$$f_1(u) = -\frac{u^{1/2}(2 - 2u + (1 + u) \log u)}{(-1 + u^{1/2})^3(1 + u^{1/2})^2},$$

$$f_2(u) = 2\frac{-1 + u^2 - 2u \log u}{(-1 + u)^3},$$

$F(u, v) =$

$$\begin{aligned} & 2\mathcal{D}_{2,2}(u, v)u^{-1}v^{1/2} + 2\mathcal{D}_{2,2}(u, v)u^{-1} + 2\mathcal{D}_{2,2}(u, v)u^{-3/2}v^{1/2} \\ & + 2\mathcal{D}_{2,2}(u, v)u^{-3/2} + 4\mathcal{D}_{3,1}(u, v)u^{-2}v^{1/2} + 4\mathcal{D}_{3,1}(u, v)u^{-2} \\ & + 4\mathcal{D}_{3,1}(u, v)u^{-5/2}v^{1/2} + 4\mathcal{D}_{3,1}(u, v)u^{-5/2} - 2\mathcal{D}_{1,2}(u, v)u^{-1/2}v^{1/2} \\ & - 2\mathcal{D}_{1,2}(u, v)u^{-1/2} - 4\mathcal{D}_{2,1}(u, v)u^{-1}v^{1/2} - 6\mathcal{D}_{2,1}(u, v)u^{-1} \\ & - 6\mathcal{D}_{2,1}(u, v)u^{-3/2}v^{1/2} - 8\mathcal{D}_{2,1}(u, v)u^{-3/2} + 2\mathcal{D}_{1,1}(u, v)u^{-1/2}v^{1/2} \\ & + 4\mathcal{D}_{1,1}(u, v)u^{-1/2} \end{aligned}$$

=

$$\begin{aligned} & (2u(-(((-1 + uv)(1 + \sqrt{u}(-1 - \sqrt{v} - (-2 + \sqrt{u} + u)v + uv^{3/2})))))/ \\ & ((-1 + \sqrt{u})(-1 + \sqrt{v}))) + (\sqrt{u}\sqrt{v}(-1 - \sqrt{u} + u + u(-2 - \sqrt{u} + 2u) \\ & \sqrt{v} + u(-1 + \sqrt{u} + u)v + u^{5/2}v^{3/2}) \log u)/((-1 + \sqrt{u})^2(1 + \sqrt{u})) \\ & + (\sqrt{v}(1 - \\ & \sqrt{u}\sqrt{v}(-1 - \sqrt{v} + v + uv(-1 + \sqrt{v} + v) + \sqrt{u}(-2 + \sqrt{v} + 2v))) \log v)/ \\ & ((-1 + \sqrt{v})^2(1 + \sqrt{v}))))/(-1 + uv)^3. \end{aligned}$$

Computations for $\partial^* k^2 \partial$ on $(1, 0)$ -forms

The symbol of $\partial^* k^2 \partial$ is equal to $c_2(\xi) + c_1(\xi)$ where

$$c_2(\xi) = \xi_1^2 k^2 + 2\tau_1 \xi_1 \xi_2 k^2 + |\tau|^2 \xi_2^2 k^2,$$

$$c_1(\xi) = (\delta_1(k^2) + \bar{\tau} \delta_2(k^2)) \xi_1 + (\tau \delta_1(k^2) + |\tau|^2 \delta_2(k^2)) \xi_2.$$

After similar computations, the second component of the scalar curvature is:

$$\begin{aligned}
& g_1(\Delta)(k^{-1}\delta_1^2(k)) + g_2(\Delta)(k^{-2}\delta_1(k)^2) \\
+ & G(\Delta_{(1)}, \Delta_{(2)})((\delta_1(k)k^{-1})(k^{-1}\delta_1(k))) \\
+ & |\tau|^2 g_1(\Delta)(k^{-1}\delta_2^2(k)) + |\tau|^2 g_2(\Delta)(k^{-2}\delta_2(k)^2) \\
+ & |\tau|^2 G(\Delta_{(1)}, \Delta_{(2)})((\delta_2(k)k^{-1})(k^{-1}\delta_2(k))) \\
+ & \tau_1 g_1(\Delta)(k^{-1}\delta_1\delta_2(k)) + \tau_1 g_2(\Delta)(k^{-2}\delta_1(k)\delta_2(k)) \\
+ & \tau_1 G(\Delta_{(1)}, \Delta_{(2)})((\delta_1(k)k^{-1})(k^{-1}\delta_2(k))) \\
+ & \tau_1 g_1(\Delta)(k^{-1}\delta_2\delta_1(k)) + \tau_1 g_2(\Delta)(k^{-2}\delta_2(k)\delta_1(k)) \\
+ & \tau_1 G(\Delta_{(1)}, \Delta_{(2)})((\delta_2(k)k^{-1})(k^{-1}\delta_1(k))) \\
+ & i\tau_2 L(\Delta_{(1)}, \Delta_{(2)})((\delta_1(k)k^{-1})(k^{-1}\delta_2(k))) \\
- & i\tau_2 L(\Delta_{(1)}, \Delta_{(2)})((\delta_2(k)k^{-1})(k^{-1}\delta_1(k)))
\end{aligned}$$

$$g_1(u) = \frac{-1 + u^2 - 2u \log u}{(-1 + u^{1/2})^3(1 + u^{1/2})^2},$$

$$g_2(u) = 2 \frac{-1 + u^2 - 2u \log u}{(-1 + u)^3},$$

$$G(u, v) =$$

$$\begin{aligned} & 2\mathcal{D}_{2,2}(u, v)u^{-1}v^{1/2} + 2\mathcal{D}_{2,2}(u, v)u^{-1} + 2\mathcal{D}_{2,2}(u, v)u^{-3/2}v^{1/2} \\ & + 2\mathcal{D}_{2,2}(u, v)u^{-3/2} + 4\mathcal{D}_{3,1}(u, v)u^{-2}v^{1/2} + 4\mathcal{D}_{3,1}(u, v)u^{-2} \\ & + 4\mathcal{D}_{3,1}(u, v)u^{-5/2}v^{1/2} + 4\mathcal{D}_{3,1}(u, v)u^{-5/2} - 4\mathcal{D}_{2,1}(u, v)u^{-1}v^{1/2} \\ & - 4\mathcal{D}_{2,1}(u, v)u^{-1} - 4\mathcal{D}_{2,1}(u, v)u^{-3/2}v^{1/2} - 4\mathcal{D}_{2,1}(u, v)u^{-3/2} \\ & - \mathcal{D}_{1,2}(u, v)v^{1/2} - \mathcal{D}_{1,2}(u, v)u^{-1/2}v^{1/2} - \mathcal{D}_{1,2}(u, v)u^{-1/2} \\ & - \mathcal{D}_{1,2}(u, v) - \mathcal{D}_{2,1}(u, v)u^{-3/2}v^{1/2} - \mathcal{D}_{2,1}(u, v)u^{-1}v^{1/2} \\ & - \mathcal{D}_{2,1}(u, v)u^{-3/2} - \mathcal{D}_{2,1}(u, v)u^{-1} - \mathcal{D}_{2,1}(u, v)u^{-1} \end{aligned}$$

$$\begin{aligned}
& -\mathcal{D}_{2,1}(u, v)u^{-1}v^{1/2} - \mathcal{D}_{2,1}(u, v)u^{-3/2} - \mathcal{D}_{2,1}(u, v)u^{-3/2}v^{1/2} \\
& + \mathcal{D}_{1,1}(u, v)u^{-1/2}v^{1/2} + \mathcal{D}_{1,1}(u, v)v^{1/2} + \mathcal{D}_{1,1}(u, v)u^{-1/2} \\
& + \mathcal{D}_{1,1}(u, v)
\end{aligned}$$

$$\begin{aligned}
= & -(\sqrt{u}(u(-1+v)^2(-1+uv(-4+u(4+v)))) \log(1/u) + (-1+u) \\
& ((1+u(-2+v))(-1+v)(-1+uv)(1+uv) + (-1+u)v \\
& (-1+u(-4+v(4+uv))) \log v)) / ((-1+\sqrt{u})^2(1+\sqrt{u})(-1+\sqrt{v})^2 \\
& (1+\sqrt{v})(-1+uv)^3),
\end{aligned}$$

$$\begin{aligned}
L(u, v) &:= -\mathcal{D}_{1,2}(u, v)u^{-1/2}v^{1/2} - \mathcal{D}_{1,2}(u, v)v^{1/2} - \mathcal{D}_{1,2}(u, v)u^{-1/2} \\
&\quad - \mathcal{D}_{1,2}(u, v) - \mathcal{D}_{2,1}(u, v)u^{-3/2}v^{1/2} - \mathcal{D}_{2,1}(u, v)u^{-1}v^{1/2} \\
&\quad - \mathcal{D}_{2,1}(u, v)u^{-3/2} - \mathcal{D}_{2,1}(u, v)u^{-1} + \mathcal{D}_{2,1}(u, v)u^{-1} \\
&\quad + \mathcal{D}_{2,1}(u, v)u^{-1}v^{1/2} + \mathcal{D}_{2,1}(u, v)u^{-3/2} + \mathcal{D}_{2,1}(u, v)u^{-3/2}v^{1/2} \\
&\quad + \mathcal{D}_{1,1}(u, v)u^{-1/2}v^{1/2} + \mathcal{D}_{1,1}(u, v)v^{1/2} + \mathcal{D}_{1,1}(u, v)u^{-1/2} \\
&\quad + \mathcal{D}_{1,1}(u, v) \\
&= (\sqrt{u}(u(-1+v)^2 \log(1/u) + (-1+u)((-1+v)(-1+uv) + (v-uv) \\
&\quad \log v)))/((-1+\sqrt{u})^2(1+\sqrt{u})(-1+\sqrt{v})^2(1+\sqrt{v})(-1+uv)).
\end{aligned}$$

Scalar curvature in terms of $\log(k)$

Lemma: For $i, j = 1, 2$, we have

$$k^{-2}\delta_i(k)\delta_j(k) = 4\frac{\Delta - \Delta^{1/2}}{\log \Delta}(\delta_i(\log k))\frac{\Delta^{1/2} - 1}{\log \Delta}(\delta_j(\log k));$$

$$k^{-1}\delta_i\delta_j(k) = 2\frac{\Delta^{1/2} - 1}{\log \Delta}(\delta_i\delta_j(\log k)) +$$

$$g(\Delta_{(1)}, \Delta_{(2)})(\delta_j(\log k)\delta_i(\log k)) + g(\Delta_{(1)}, \Delta_{(2)})(\delta_i(\log k)\delta_j(\log k)),$$

where

$$g(u, v) := 4\frac{(\sqrt{uv} - 1)\log u - (\sqrt{u} - 1)\log(uv)}{\log v \log u \log(uv)},$$

and $\Delta_{(i)}$ signifies the action of Δ on the i -th factor of the product.

Final formula for the scalar curvature (Connes-Moscovici, Fathizadeh-K.)

Theorem: The scalar curvature of (A_θ, τ, k) , up to an overall factor of $\frac{-\pi}{\tau_2}$, is equal to

$$\begin{aligned} & R_1(\log \Delta)(\Delta_0(\log k)) + \\ & R_2(\log \Delta_{(1)}, \log \Delta_{(2)}) \left(\delta_1(\log k)^2 + |\tau|^2 \delta_2(\log k)^2 + \tau_1 \{ \delta_1(\log k), \delta_2(\log k) \} \right) + \\ & iW(\log \Delta_{(1)}, \log \Delta_{(2)}) \left(\tau_2 [\delta_1(\log k), \delta_2(\log k)] \right) \end{aligned}$$

where

$$R_1(x) = -\frac{\frac{1}{2} - \frac{\sinh(x/2)}{x}}{\sinh^2(x/4)},$$

$$R_2(s, t) = (1 + \cosh((s + t)/2)) \times \frac{-t(s + t) \cosh s + s(s + t) \cosh t - (s - t)(s + t + \sinh s + \sinh t - \sinh(s + t))}{st(s + t) \sinh(s/2) \sinh(t/2) \sinh^2((s + t)/2)},$$

$$W(s, t) = -\frac{(-s - t + t \cosh s + s \cosh t + \sinh s + \sinh t - \sinh(s + t))}{st \sinh(s/2) \sinh(t/2) \sinh((s + t)/2)}.$$

The limiting case

In the commutative case, the above modular curvature reduces to a constant multiple of the [formula of Gauss](#):

$$\frac{1}{\tau_2} \delta_1^2(\log k) + \frac{|\tau|^2}{\tau_2} \delta_2^2(\log k) + 2 \frac{\tau_1}{\tau_2} \delta_1 \delta_2(\log k).$$

First application: the Gauss-Bonnet theorem for A_θ

Spectral formulation of the Gauss-Bonnet theorem:

$$\zeta(0) + 1 = \frac{1}{12\pi} \int_{\Sigma} R \operatorname{vol}(g) = \frac{1}{6} \chi(\Sigma)$$

Theorem (Connes-Tretkoff; Fathizadeh-K.): Let $\theta \in \mathbb{R}$, $\tau \in \mathbb{C} \setminus \mathbb{R}$, $k \in A_\theta^\infty$ be a positive invertible element. Then

$$\operatorname{Trace}(\Delta^{-s})|_{s=0} + 2 = \mathfrak{t}(R) = 0,$$

where Δ is the Laplacian and R is the scalar curvature of the spectral triple attached to (A_θ, τ, k) .

Second application: conformal anomaly (Connes-Moscovici)

Polyakov's conformal anomaly formula:

$$\log \det(\Delta) = -\zeta'_{\Delta}(0)$$

$$\log \det(\Delta) - \log \det(\Delta_0) = \log \varphi_0(e^{-h}) + \varphi_0(\tilde{R}).$$

Scalar curvature for higher NC tori

- ▶ Let $\theta : V \otimes V \rightarrow \mathbb{R}$ alternating form, $\Lambda \subset V$ a cocompact lattice. $A_\theta = A_\theta(V, \Lambda)$ by generators and relations:

$$U_\alpha U_\beta = e^{\pi i \theta(\alpha, \beta)} U_{\alpha + \beta}, \quad \text{for all } \alpha, \beta \in \Lambda.$$

- ▶ The dual torus V/Λ^* acts on A_θ by

$$\lambda_s(U_\alpha) = e^{2\pi i \langle s, \alpha \rangle} U_\alpha.$$

- ▶ The differential of the action λ_s at identity defines a Lie algebra map

$$\delta : V \rightarrow \text{Der}(A_\theta^\infty, A_\theta^\infty)$$

where $A_\theta^\infty =$ space of smooth vectors of λ_s . It is a dense *-subalgebra of A_θ .

- ▶ Let $\mathfrak{t} : A_\theta \rightarrow \mathbb{C}$ denote the canonical tracial state. Its restriction on smooth elements is given by

$$\mathfrak{t}\left(\sum_{\alpha \in \mathbb{Z}^n} a_\alpha U_\alpha\right) = a_0.$$

Noncommutative complex tori

- ▶ Let $\Lambda \subset V$ be a lattice in a complex v.s.; e_1, \dots, e_n a basis of V , and $\lambda_1, \dots, \lambda_{2n}$ a basis of Λ . Express $\lambda_1, \dots, \lambda_{2n}$ in terms of e_1, \dots, e_n and obtain an n by $2n$ matrix $\mathcal{M} = (A, B)$ with $A, B \in M_n(\mathbb{C})$ with

$$\lambda_j = \sum_{i=1}^n \mathcal{M}_{ij} e_i, \quad j = 1, \dots, 2n$$

- ▶ Let dz_1, \dots, dz_n denote the basis of $V^* = \text{Hom}_{\mathbb{C}}(V, \mathbb{C})$, dual to e_1, \dots, e_n , and let dx_1, \dots, dx_{2n} denote the basis of $V_{\mathbb{R}}^* = \text{Hom}_{\mathbb{R}}(V, \mathbb{R})$, dual to $\lambda_1, \dots, \lambda_{2n}$.

So that

$$dz_i = \sum_{j=1}^{2n} \mathcal{M}_{ij} dx_j, \quad d\bar{z}_i = \sum_{j=1}^{2n} \bar{\mathcal{M}}_{ij} dx_j$$

We get the differential operators

$$\begin{aligned} \delta_i = \delta_{\lambda_i} & : A_\theta^\infty \rightarrow A_\theta^\infty, & i = 1, \dots, 2n \\ \partial_i = \delta_{e_i} & : A_\theta^\infty \rightarrow A_\theta^\infty, \end{aligned}$$

These derivations satisfy the relations

$$\delta_j = \sum_{i=1}^n \mathcal{M}_{ij} \partial_i, \quad 1 \leq j \leq n$$

Dolbeault complex

Let

$$\Omega^{p,q} := A_{\theta}^{\infty} \otimes \wedge^p V_{(1,0)}^* \otimes \wedge^q V_{(0,1)}^*$$

$$\partial_i : \Omega^{p,q} \rightarrow \Omega^{p+1,q}, \quad \bar{\partial}_i : \Omega^{p,q} \rightarrow \Omega^{p,q+1}.$$

$$\partial_i(adz_I \wedge d\bar{z}_J) = \sum_i \partial_i(a) dz_i \wedge dz_I \wedge d\bar{z}_J$$

$$\bar{\partial}_i(adz_I \wedge d\bar{z}_J) = \sum_i \bar{\partial}_i(a) d\bar{z}_i \wedge dz_I \wedge d\bar{z}_J$$

We have:

$$\partial^2 = 0, \quad \bar{\partial}^2 = 0, \quad \partial\bar{\partial} + \bar{\partial}\partial = 0.$$

Noncommutative abelian varieties

- ▶ Fix $\Omega \in \mathfrak{H}_n \subset M_n(\mathbb{C})$, *Siegel upper half space*:

$$\Omega^t = \Omega \quad \text{and} \quad \text{Im } \Omega > 0$$

- ▶ A complex torus V/Λ is projective (hence Kaehler) iff there exists a basis (e_1, \dots, e_n) of V , and a basis $(\lambda_1, \dots, \lambda_{2n})$ of Λ such that

$$(\lambda_1, \dots, \lambda_{2n}) = (\Delta, \Omega)$$

where $\Delta = \text{diag}(k_1, \dots, k_n)$, $k_i \in \mathbb{Z}$, is a diagonal matrix, and $\Omega \in \mathfrak{H}_n$.

Dolbeault Laplacian v.s. de Rham Laplacian

If M is Kaehler then

$$\Delta_{de\ Rham} = 2\Delta_{Dolbeault}$$

So on Kaehler Manifolds Dolbeault Laplacian and de Rham Laplacian have same spectrum and same spectral zeta functions.

- ▶ Any Riemann surface is Kaehler.
- ▶ 4-tori with perturbed metric are not Kaehler.

- ▶ So we have to work directly with de Rham Laplacian.

Laplacian for $\Omega = \sqrt{-1} I_2$

In this case the de Rham differential can be written as

$$d = \partial + \bar{\partial} : A_\theta \rightarrow A_\theta^4,$$

where,

$$\partial_1 = \frac{1}{2}(\delta_1 - i\delta_3), \quad \partial_2 = \frac{1}{2}(\delta_2 - i\delta_4),$$

$$\bar{\partial}_1 = \frac{1}{2}(\delta_1 + i\delta_3), \quad \bar{\partial}_2 = \frac{1}{2}(\delta_2 + i\delta_4).$$

Flat Laplacian

$$d^*d = \partial_1^*\partial_1 + \partial_2^*\partial_2 + (\bar{\partial}_1)^*\bar{\partial}_1 + (\bar{\partial}_2)^*\bar{\partial}_2$$

Perturbing the metric

Perturb the volume form on A_θ by e^{-2h} , $h \in A_\theta$. We have to consider the factor e^{-h} for 1-forms.

$$(x, y)_0 = t(y^* x e^{-2h}), \quad H_0 = \overline{A_\theta}^{(\cdot, \cdot)_0}$$

$$(x, y)_1 = t(y^* x e^{-h}), \quad H_1 = \overline{A_\theta}^{(\cdot, \cdot)_1}$$

Lemma: The perturbed Laplacian $d^*d : H_0 \rightarrow H_0$ is anti-unitarily equivalent to $\Delta' =$

$$e^h \bar{\partial}_1 e^{-h} \partial_1 e^h + e^h \partial_1 e^{-h} \bar{\partial}_1 e^h + e^h \bar{\partial}_2 e^{-h} \partial_2 e^h + e^h \partial_2 e^{-h} \bar{\partial}_2 e^h.$$

Symbol of the perturbed Laplacian

With $k = e^{h/2}$, the symbol of Δ' is

$$a_2(\xi) + a_1(\xi) + a_0(\xi) : \mathbb{R}^4 \rightarrow A_\theta^4,$$

where

$$a_2(\xi) = \frac{1}{2}k^2 \sum_{i=1}^4 \xi_i^2, \quad a_1(\xi) = \frac{1}{2} \sum_{i=1}^4 \delta_i(k^2)\xi_i,$$

$$a_0(\xi) = \frac{1}{2} \sum_{i=1}^4 (\delta_i^2(k^2) - \delta_i(k^2)k^{-2}\delta_i(k^2)).$$

Approximate the resolvent $(\Delta' - \lambda)^{-1}$ by a pseudodifferential operator B_λ whose symbol has an expansion of the form

$$b_0(\xi, \lambda) + b_1(\xi, \lambda) + b_2(\xi, \lambda) + \dots$$

where $b_j(\xi, \lambda)$ is a symbol of order $-2 - j$.

Computing $b_2(\xi, \lambda)$

$$b_2 = - \sum_{\substack{2+j+\ell_1+\ell_2+\ell_3+\ell_4-k=2, \\ 0 \leq j < 2, 0 \leq k \leq 2}} \frac{1}{\ell!} \partial^\ell (b_j) \delta^\ell (a_k) b_0.$$

$$b_2(\xi, \lambda) =$$

$$\begin{aligned} & -\frac{1}{2} b_0 k \delta_1 [\delta_1[k]] b_0 - \frac{1}{2} b_0 k \delta_2 [\delta_2[k]] b_0 - \frac{1}{2} b_0 k \delta_3 [\delta_3[k]] b_0 - \\ & \frac{1}{2} b_0 k \delta_4 [\delta_4[k]] b_0 - \frac{1}{2} b_0 \delta_1[k] \delta_1[k] b_0 - \frac{1}{2} b_0 \delta_1 [\delta_1[k]] b_0 k - \\ & \frac{1}{2} b_0 \delta_2[k] \delta_2[k] b_0 - \frac{1}{2} b_0 \delta_2 [\delta_2[k]] b_0 k - \frac{1}{2} b_0 \delta_3[k] \delta_3[k] b_0 - \\ & \frac{1}{2} b_0 k \delta_1[k] \frac{1}{k} \delta_1[k] b_0 + (\text{more than 2200 similar terms}) \dots \end{aligned}$$

Homogeneity argument for contour integrals

$$R = \frac{1}{2\pi i} \int_{\mathbb{R}^4} \int_C e^{-\lambda} b_2(\xi, \lambda) d\lambda d\xi$$

Integration by parts for inner integral \implies

$$\int_C b_2(\xi, \lambda) d(-e^{-\lambda}) = \int_C e^{-\lambda} D_\lambda(b_2(\xi, \lambda)) d\lambda$$

Define $\beta(\lambda) := \int_{\mathbb{R}^4} D_\lambda(b_2(\xi, \lambda)) d\lambda d\xi$. $\beta(\lambda)$ is homogenous of order -1 .

$$\beta(t\lambda) = t^{-1}\beta(\lambda).$$

So we have

$$R = \frac{1}{2\pi i} \int_C e^{-\lambda} \beta(\lambda) d\lambda = -\beta(-1) = \frac{-1}{2\pi i} \int_{\mathbb{R}^4} D_\lambda(b_2(\xi, -1)) d\xi.$$

Scalar curvature of NC 4-Tori (Fathizadeh-Ghorbanpour-k)

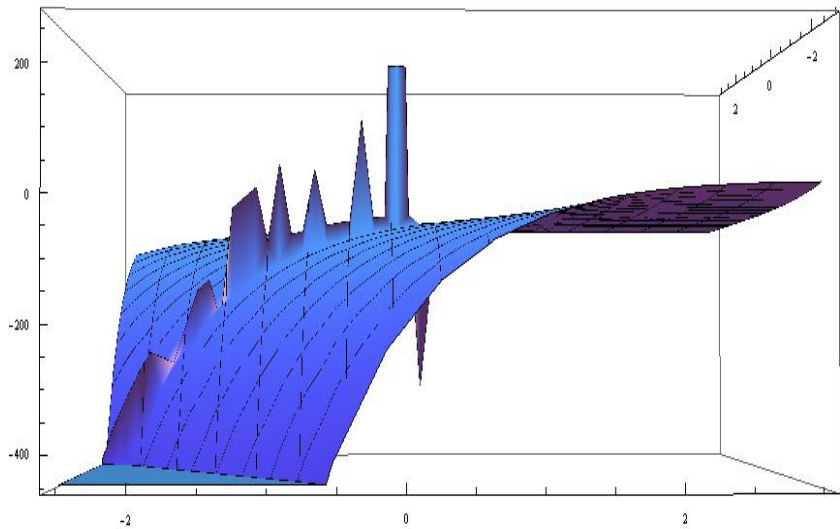
$$R = \frac{1}{2}e^{-h}K(\nabla)\left(\sum_{i=1}^4 \delta_i^2(h)\right) + \frac{1}{4}e^{-h}H(\nabla_{(1)}, \nabla_{(2)})\left(\sum_{i=1}^4 \delta_i(h)^2\right)$$

where $\Delta(x) = k^{-2}xk^2 = e^{-h}xe^h$, $\nabla(x) = \log \Delta(x) = [-h, x]$ for $x \in A_\theta$,

$$K(s) = \frac{2\pi^2 e^{-s} (e^s - 1)}{s},$$

and

$$H(s, t) = \frac{8\pi^2 e^{-3s/2-t}}{st} \left(-\frac{e^{s/2}(e^{\frac{s}{2}+\frac{t}{2}} + 1)(\frac{s}{2} - e^{t/2}(-\frac{1}{2}e^{s/2}t + \frac{s}{2} + \frac{t}{2}))}{\frac{s}{2} + \frac{t}{2}} \right. \\ \left. + e^{s+\frac{t}{2}} \sinh(\frac{s}{2}) \sinh(\frac{t}{2}) - 2 \sinh(\frac{s}{4}) \sinh(\frac{t}{4})((\frac{s}{2} + \frac{t}{2}) \coth(\frac{s}{2} + \frac{t}{2}) + \frac{s}{2} + \frac{t}{2} - 1) \right)$$



Two variable function

A check: the classical limit

We have

$$\lim_{s \rightarrow 0} K(s) = 2\pi^2, \quad \lim_{s, t \rightarrow 0} H(s, t) = -2\pi^2.$$

Therefore, in the commutative case, the above scalar curvature reduces to

$$\pi^2 e^{-h} \sum_{i=1}^4 (\delta_i^2(h) - \frac{1}{2} \delta_i(h)^2),$$

which, up to a constant multiple, is precisely the scalar curvature of the ordinary 4-torus when the standard metric is perturbed by a factor of e^{-h} .